## LA-UR-22-30536

#### Approved for public release; distribution is unlimited.

 Title:
 Adding Delta Tracking to the MCNP Code

Author(s): Josey, Colin James

Intended for: 2022 MCNP User Symposium, 2022-10-17/2022-10-21 (Los Alamos, New Mexico, United States)

**Issued:** 2022-10-10









Los Alamos National Laboratory, an affirmative action/equal opportunity employer, is operated by Triad National Security, LLC for the National Nuclear Security Administration of U.S. Department of Energy under contract 89233218CNA000001. By approving this article, the publisher recognizes that the U.S. Government retains nonexclusive, royalty-free license to publish or reproduce the published form of this contribution, or to allow others to do so, for U.S. Government purposes. Los Alamos National Laboratory requests that the publisher identify this article as work performed under the auspices of the U.S. Department of Energy. Los Alamos National Laboratory strongly supports academic freedom and a researcher's right to publish; as an institution, however, the Laboratory does not endorse the viewpoint of a publication or guarantee its technical correctness.



# Adding Delta Tracking to the MCNP Code

**Colin Josey** 

XCP-3 Monte Carlo Codes

MCNP User Symposium 18 October 2022



#### Outline

#### Theory

**MCNP** Implementation

#### Testing



Section 1

## Theory



#### **Background & Motivation**

Within the MCNP code, particles are tracked using "surface-to-surface" tracking:

- Break geometry into regions with constant materials.
- Sample distance to collision  $\delta_c$ , compute distance to surface,  $\delta_s$
- If  $\delta_c > \delta_s$ , move to surface
- Else, move to collision and sample collision

Implementation is easy, as constant materials means constant collision probability  $(\Sigma_t)$ , and the distance sampling is analytic:

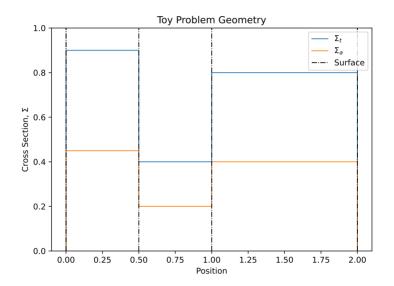
$$\delta_c = \frac{-\ln(\xi)}{\Sigma_t}$$

Similarly, constant materials means tallies have a constant multiplier over each track, and one can use track-length tallies (which are used extensively within the MCNP code).



#### **Toy Problem**

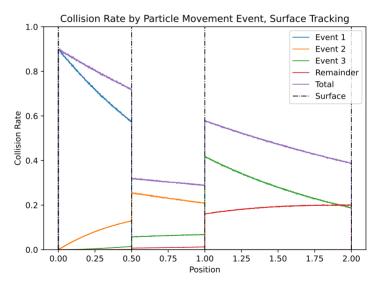
Let us consider a toy problem for demonstrating algorithms.





#### **Surface Tracking**

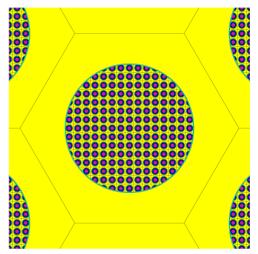
Event: any iteration of the transport loop (here, collision or surface crossing), requiring data lookups and resampling.





#### **TRISO Problems**

- TRISO fuel is very small
- A full reactor can have well in excess of 5 billion surfaces
- Every crossing requires many computations
- This results in poor performance



TRISO Fuel Element (cylinder is 1.5 cm across)



#### **Alternatives**

- Numerical Methods
  - Can handle continuously varying materials well
  - Discontinuities require either stopping the particle or taking many steps near the discontinuity
- Delta Tracking
  - Focus of this talk
  - Can handle discontinuous materials
  - Used in the Serpent Monte Carlo code



#### The Delta Tracking Method

- 1. Add to the collision probability  $\Sigma_t$  a probability of not colliding  $\Sigma_{nc}$ .
- 2. Select  $\Sigma_{nc}$  such that  $\Sigma_t + \Sigma_{nc}$  is constant and  $\Sigma_{nc}$  is non-negative.  $\Sigma_t + \Sigma_{nc}$  is now called the "Majorant",  $\Sigma_m$ .
- 3. Track with this new cross section. Distance is now:

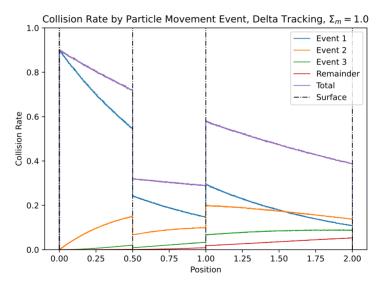
$$\delta_c = \frac{-\ln(\xi_1)}{\Sigma_m}$$

- 4. Move particle  $\delta_c$ .
- 5. Randomly sample if a collision was real or not:
- 5.1 If  $\xi_2 < \frac{\Sigma_t}{\Sigma_m}$ , collide
- 5.2 Else, leave particle as-is.



#### **Delta Tracking**

Using delta tracking, one notes that every part of the geometry is tallied in the first event, and fewer events are needed to make it through on average.





#### Advantages and Disadvantages to Delta Tracking

Advantages:

 Adding more surfaces does not affect performance (aside from possibly increasing cell-at-point-of-collision lookup time)

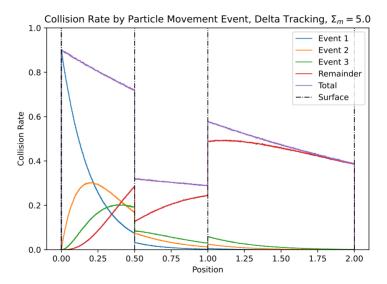
Disadvantages:

- A large majorant will cause more collisions, harming performance
  - Burnable absorbers are key culprit
- Cannot use track-length tally estimators for most values of interest without issue
  - Exception is flux tallies in a structured mesh, where distance in cell is easy to compute and there is no discontinuous weighting function



## Large Majorant Example

If the majorant is increased (say, due to burnable absorbers), events occur more often, reducing efficiency.





#### **Tallies**

As mentioned earlier, delta tracking does not support track-length estimators in most configurations:

- For cell tallies, would need to determine if each track entered a cell and for how long
- ► For reaction rate tallies, would need the same for materials

Computing this information would nullify the delta tracking advantage. So, instead, at every event, we tally:

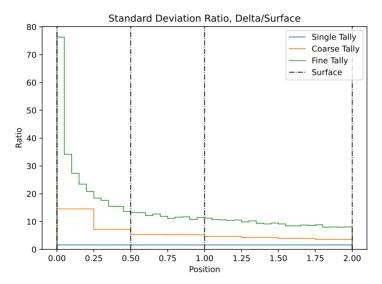
$$\mathsf{score} = \frac{\mathsf{weight} \times \mathsf{multiplier} \text{ at event}}{\Sigma_m}$$

(Pseudo-)collision tallies tend to have a larger variance than track-length tallies, particularly as tally domain size shrinks.



## **Tally Comparisons**

As the tallies get finer, track-length tallies statistically outperform collision tallies. Ratios are exaggerated by being 1D.





#### **Figure of Merit**

While the variance of the tallies may be higher, if the simulation runs faster, you can run more particles to make up for it.

variance 
$$\propto \! \frac{1}{\text{number of simulated particles}} \propto \frac{1}{\text{time}}$$

Based on the above relationship, (variance  $\times$  time) should be roughly constant for any given algorithm. So, the figure of merit for Monte Carlo, FOM is given by:

$$\mathsf{FOM} = \frac{1}{\sigma^2 t}$$

The ratio of the FOM between algorithms shows, roughly, how much more efficient one algorithm is compared to another.



#### Section 2

**MCNP** Implementation



#### **Computing the Majorant**

There are several ways one can compute the majorant:

- 1. Get the global maximum  $\Sigma_t$
- 2. Create a function such that  $\Sigma_m(E) \ge \Sigma_t(E)$  for all positions

This latter form is made possible by the fact that neutrons and photons do not change energy except during collision (in the physics implemented in MCNP).

Making it a function of energy is beneficial. As shown earlier, a large majorant reduces efficiency. This form will reduce  $\Sigma_m$  over most energies.



#### **Data Components**

For each particle, the majorant must take into consideration all contributions. In the MCNP code, these are:

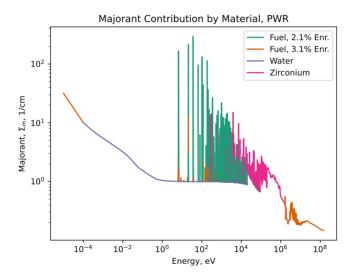
- Neutrons
  - Basic reaction rates linear-linear representation
  - Unresolved resonances several possible representations
  - ►  $S(\alpha, \beta)$  two components are linear-linear, one is piecewise functions proportional to 1/E
- Photons
  - Photoatomic log-log representation
  - Photonuclear linear-linear representation

 $\Sigma_m(E)$  was chosen to be linear-linear interpolated, as it makes several datasets trivial and the rest can be quickly approximated as linear-linear.



#### Majorant

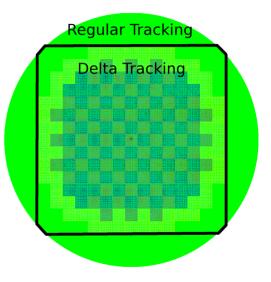
The majorant is then computed during simulation load:





#### **Transport Modifications**

- A hybrid approach was used:
  - If in a "delta-tracking universe", use the delta tracking code path
  - If not, use the regular surface-to-surface code path
- This allows one to get performance benefits inside of a core, while still allowing surface tallies in the perimeter.





#### **Gaps in Implementation**

The MCNP code has decades of features implemented assuming surface-to-surface tracking. To support everything in the first pass would be folly.

The following are known to not work with delta tracking at this time:

- Most variance reduction many need rederiving/reconsidering
- Unstructured/structured mesh geometry uses a completely different code path
- All non-neutron/photon particles most of the benefits are invalid for charged particles, for example

Other particles are still allowed, they just ignore the delta tracking flag.



Section 3

#### Testing



## Testing

With such a radical change in the transport algorithm, care needed to be taken to ensure things did not break.

Generated two *k*-eigenvalue test problems:

- 1. A 2D PWR
- 2. A simplified TRISO reactor

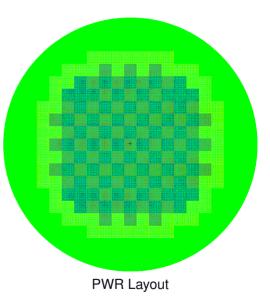
Configuration:

- Each run 50 times, with and without delta tracking
- Neutrons, photons, and photonuclear physics was enabled
- Cell tallies performed in several cells
- FMESH energy deposition tallies performed over whole core



#### 2D PWR

- 1,000,000 particles per generation
- 100 inactive generations
- 200 active generations
- ENDF/B-VIII.0 data
- ▶  $15 \times 15$  FMESH tallies
- Tallies in fuel, clad, water





#### **2D PWR Tallies**

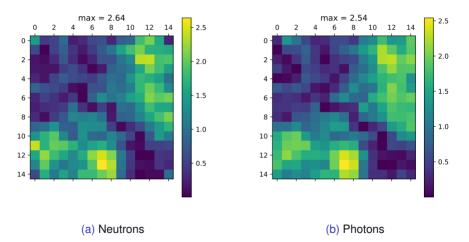
Value	$\delta$ , in std-dev	FOM Ratio
k-eigenvalue	0.50	1.65
Neutron $\Sigma_a$ , fuel	0.45	1.34
Neutron $\Sigma_a$ , clad	1.22	0.70
Neutron $\Sigma_a$ , water	0.04	1.28
Photon $\Sigma_t$ , fuel	0.74	1.54
Photon $\Sigma_t$ , clad	0.94	1.14
Photon $\Sigma_t$ , water	1.37	1.67

In all cases, the difference in units of standard deviations does not indicate any issues with the implementation. Further, the maximum relative error between the two transport modes was  $5.4 \times 10^{-5}$ , putting a tight bound on any error hidden by statistics.

Performance wise, the figure of merit increased a moderate amount for most tallies. The cladding, being quite thin, falls afoul of the increased variance for small regions issue.



#### **2D PWR FMESH Accuracy**



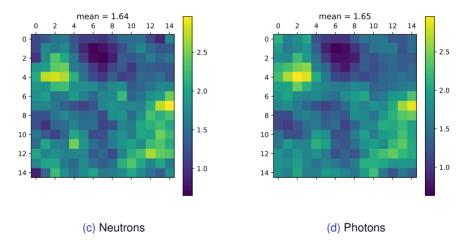
Difference between tallies, in units of standard deviations

Both FMESH tally results show no unexpected discrepancies.



#### **2D PWR FMESH FOM**



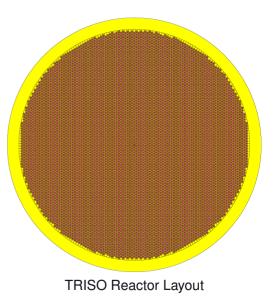


Delta tracking improved performance moderately, but it was spatially non-uniform.



#### **TRISO Model**

- 100,000 particles per generation
- 100 inactive generations
- 200 active generations
- ENDF/B-VIII.0 data
- ▶  $22 \times 22$  FMESH tallies
- Tallies in fuel, TRISO layers





#### **TRISO Model Tallies**

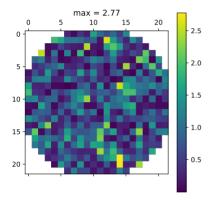
Value	$\delta$ , in std-dev	FOM Ratio
k-eigenvalue	0.41	13.49
Neutron $\Sigma_a$ , fuel	1.42	13.39
Neutron $\Sigma_a$ , layer 1	0.71	3.77
Neutron $\Sigma_a$ , layer 2	0.86	0.86
Neutron $\Sigma_a$ , layer 3	0.03	4.58
Neutron $\Sigma_a$ , layer 4	0.13	1.63
Photon $\Sigma_t$ , fuel	0.30	19.78
Photon $\Sigma_t$ , layer 1	0.80	18.78
Photon $\Sigma_t$ , layer 2	1.03	8.55
Photon $\Sigma_t$ , layer 3	0.25	8.75
Photon $\Sigma_t$ , layer 4	0.58	8.73

No indication of errors in the implementation. The maximum relative error was  $4.8 \times 10^{-4}$ . FOM improvements are much larger than the PWR results for most of the geometry.



#### **TRISO FMESH Accuracy**

Difference between tallies, in units of standard deviations

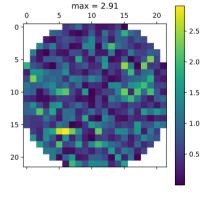


(e) Neutrons

(f) Photons

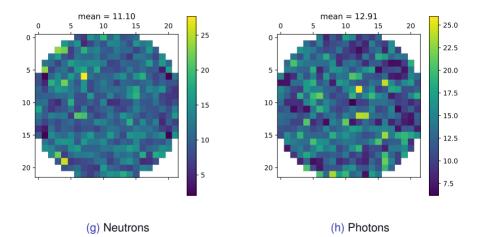
Still no indication of abnormal behavior.





#### **TRISO FMESH FOM**

Ratio of figure of merit, delta over surface-to-surface



Performance improvements are more uniform and of larger magnitude than the PWR simulation.



#### Summary

- Small improvements in regular geometries like a PWR
- Large performance improvements in fine geometries like TRISO reactors
- Testing showed no indication of implementation error in any test performed
- Further testing and cross-feature checking will be done for a future production release of the MCNP code.



#### **Acknowledgements**

This work was supported by the Los Alamos National Laboratory LDRD Program.



