

05520

LA-UR-09-?????

Approved for public release;
distribution is unlimited.

Title: Recent and Future MCNP6 Capabilities for Radiation Effects / Hostile Environments (U)

Author(s): Roger Martz

Intended for: LLNL Meeting, September 9, 2009.



Los Alamos National Laboratory, an affirmative action/equal opportunity employer, is operated by the Los Alamos National Security, LLC for the National Nuclear Security Administration of the U.S. Department of Energy under contract DE-AC52-06NA25396. By acceptance of this article, the publisher recognizes that the U.S. Government retains a nonexclusive, royalty-free license to publish or reproduce the published form of this contribution, or to allow others to do so, for U.S. Government purposes. Los Alamos National Laboratory requests that the publisher identify this article as work performed under the auspices of the U.S. Department of Energy. Los Alamos National Laboratory strongly supports academic freedom and a researcher's right to publish; as an institution, however, the Laboratory does not endorse the viewpoint of a publication or guarantee its technical correctness.

Recent and Future MCNP6 Capabilities for Radiation Effects / Hostile Environments

ABSTRACT

LA-UR-09-00???

This presentation provides an overview of some recent and future MCNP6 capabilities that may be useful for radiation effects / hostile environment studies.

Outline

- **Nested DXTRAN**
 - complete in MCNP6 trunk
- **Variable Density (for atmospheric studies)**
 - currently being implemented and tested in MCNP6
- **Deterministic adjoint generated weight windows**
 - some work already completed in MCNP6
 - some funding for additional work in FY10
 - work to do beyond FY10

Nested DXTRAN

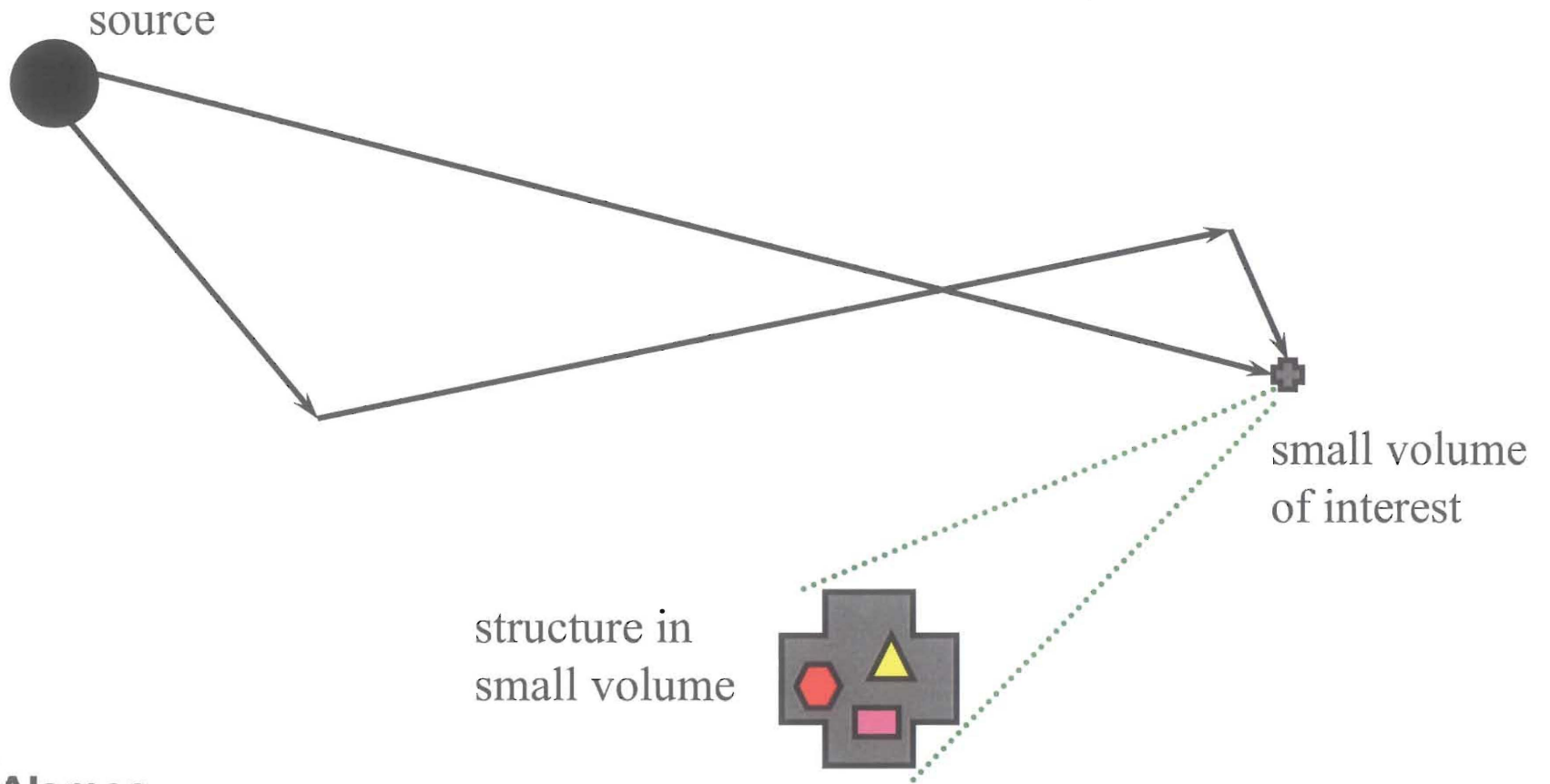
A Variance Reduction Technique

Motivation

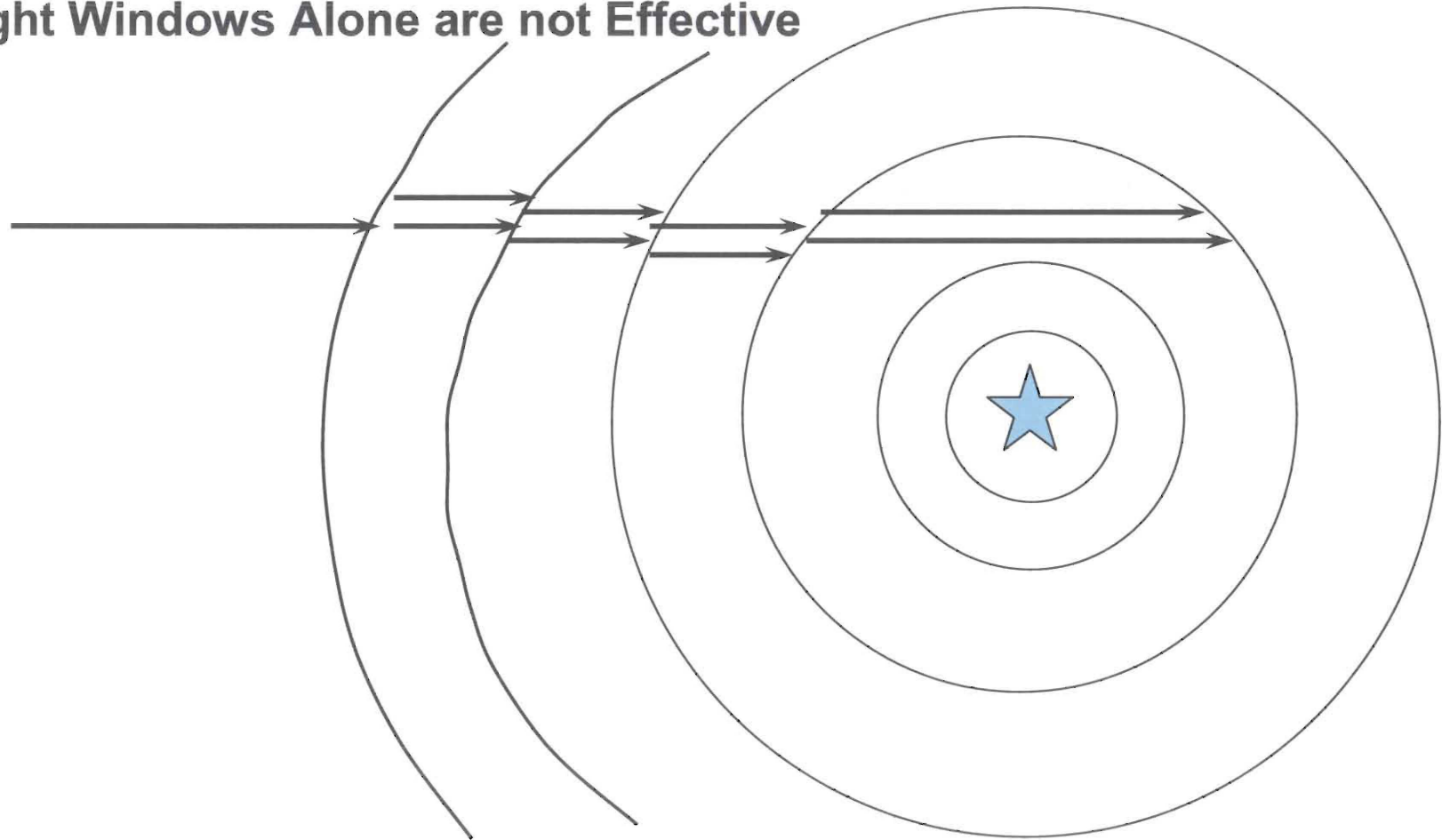
- **Current DXTRAN technique has problems in scattering media.**
- **Nested DXTRAN spheres (combined with weight windows) ensures appropriate weights in and about the tally region.**
- **Helps solve many extremely angle-dependent problems such as computing fluxes in small bodies located in the atmosphere.**

Nested DXTRAN

- Particle random walks are very unlikely to get to a small volume



Weight Windows Alone are not Effective



Splitting does not solve the angle problem

Analog vs Biased Densities / Weighting

LA-UR-09-00???

DXTRAN partitions the random walk into two parts:

1. Non-DXTRAN particle needs no weight correction - sampled same as w/o DXTRAN.
2. The DXTRAN particle needs a weight correction because it is sampled from a biased density.

$$w_{dxtran} = w \frac{p(\Omega)}{b(\Omega)} e^{-S(\Omega)}$$

biased DXTRAN exit angle pdf

true exit angle pdf

$$b(\Omega) = \frac{1}{\int_{\Omega_sphere} d\Omega}$$

The integration is over all directions pointed toward the DXTRAN sphere from the source or collision exit point.

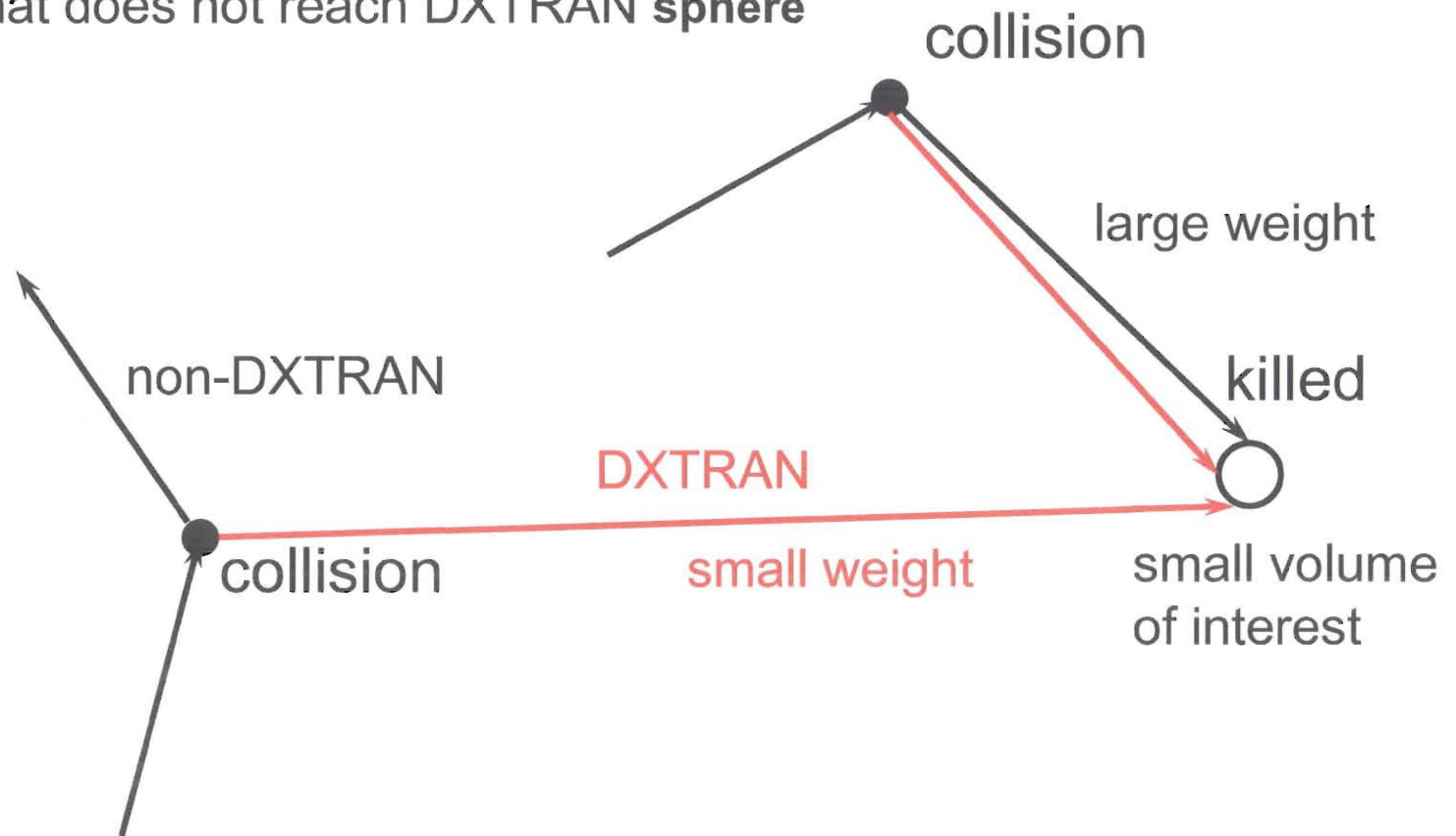
A fair game because the expected weight is preserved.

Current DXTRAN is a Partial Solution

LA-UR-09-00???

Split particle upon exiting collision:

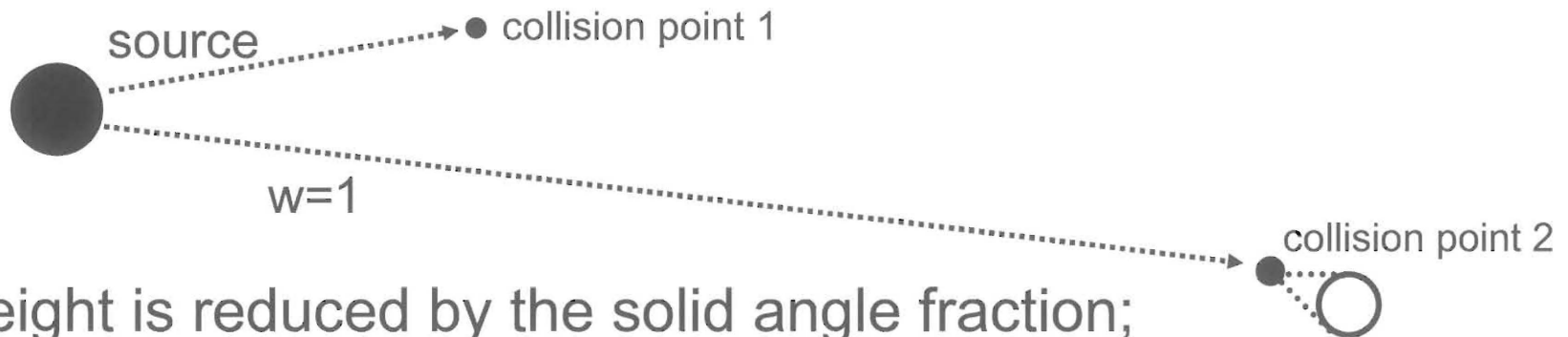
1. weight that reaches DXTRAN sphere
2. weight that does not reach DXTRAN sphere



Current DXTRAN is *Only* a Partial Solution

LA-UR-09-00???

It is still possible to get large weights in the tally region.



Weight is reduced by the solid angle fraction; there is not much reduction if the collision is close to the sphere.

$$w_{dxtran} = w \frac{p(\Omega)}{b(\Omega)} e^{-S(\Omega)} \quad b(\Omega) = \frac{1}{\int_{\Omega_sphere} d\Omega}$$

Note that b will be very large for point 1 but can be as small as

$$b(\Omega) = \frac{1}{2\pi} \quad \text{for point 2 near the sphere surface.}$$

Nesting Precludes Large Weights

- The first (inner) DXTRAN sphere prohibits a full weight particle from entering the first sphere, but the full weight particle might still collide just before the first sphere (point 2).
- A second (outer) DXTRAN sphere surrounding the first ensures that the particle colliding at point 2 cannot have full weight because its weight was reduced by the DXTRAN biasing to the second sphere.

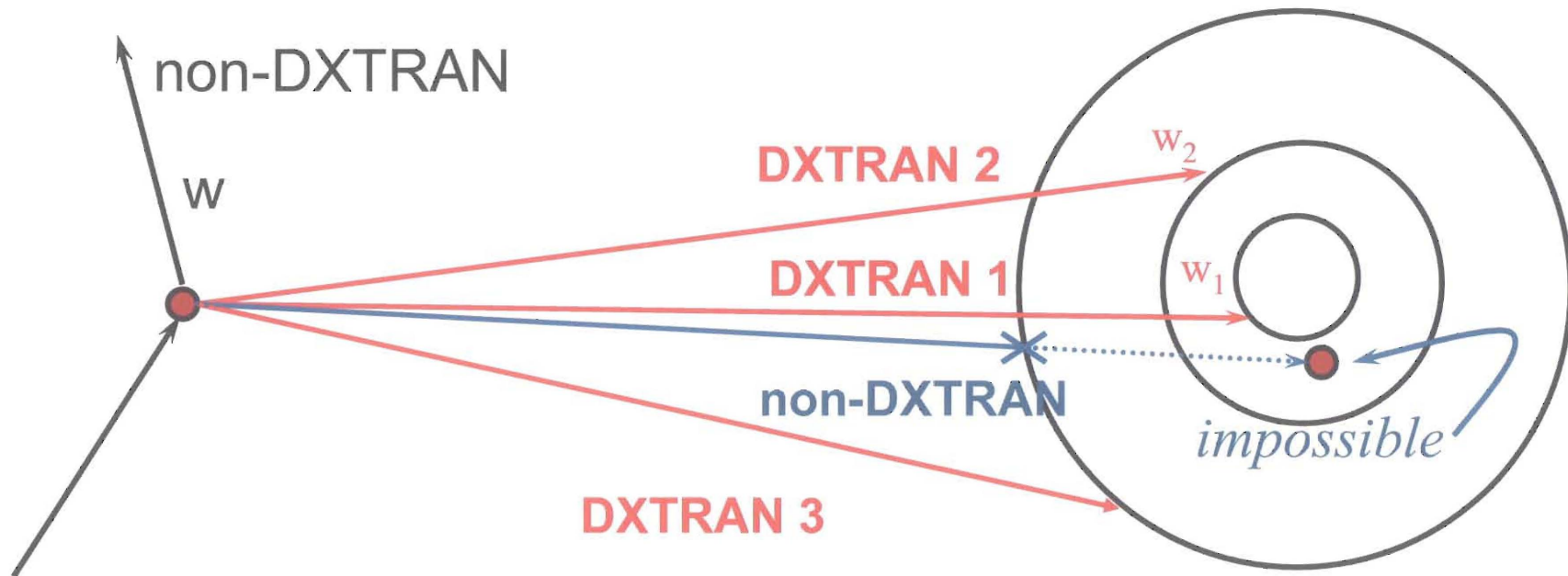


$$w_{dxtran} = w \frac{p(\Omega)}{b(\Omega)} e^{-S(\Omega)}$$

Nesting Solves Weight Fluctuations

LA-UR-09-00???

- Sphere n shields all spheres $k < n$ from large weights.
- Continue nesting DXTRAN spheres until the weight window by itself can get enough particles into the vicinity of the outermost DXTRAN sphere with appropriate weight.



Variable Density Model Improvement

Motivation

- Improve overall accuracy in atmospheric radiation transport problems.
- Generalize variable density treatment for other potential applications.

■ sub-stepping

X-5 Monte Carlo Team, MCNP – A General Monte Carlo N-Particle Transport Code, Version 5, LA-UR-03-1987, Los Alamos, 2003.

■ delta-tracking, hole-tracking, self-scattering, pseudo-scattering, Woodcock tracking

E. R. Woodcock et al., Proc. CACMRP65, ANL-7050, ANL, 1965;
C. J. Everett, E. D. Cashwell, and R. G. Schrandt, LA-5089-MS, US-34, 1972;
L. L. Carter, E. D. Cashwell, and W. M. Taylor, Nucl. Sci. Eng. 48 (1972) 40;
S. N. Cramer, ORNL Report ORNL/TM-48880, Oak Ridge, 1977.

■ direct method

F. B. Brown and W. R. Martin, Proc. NMCS03, LA-UR-02-6530, 2002;
F. B. Brown and W. R. Martin, Proc. NECDC 2002, LA-UR-02-6567, 2002;
F. Brown, D. Griesheimer, W. Martin, Proc. PHYSOR-2004, LA-UR-04-0732.

■ mass integral scaling

The Chosen Approach: Mass Integral Scaling

LA-UR-09-00???

In an infinite homogeneous medium with an isotropic point source, the $4\pi R^2$ fluence (time integrated flux) or dose is a function only of ρR , the mass per unit area between the source and receiver, known as “mass range” or “areal density”.

Mathematically, the mass scaling law can be derived:

A) using the Boltzmann transport equation:

C. D. Zerby, “Radiation Flux Transformation as a Function of Density in an Infinite Medium with Anisotropic Point Sources,” ORNL Report 2100, Oak Ridge, 1965.

B) from the diffusion equation:

Raymond A. Shulstad, “An Evaluation of Mass Integral Scaling as Applied to the Atmospheric Radiation Transport Problem,” PhD thesis, Air Force Institute of Technology (AFIT/SA), Wright-Patterson AFB, OH, DS/PH/76-3, 1976.

Previous Implementations

LA-UR-09-00???

Mass Integral Scaling Approach has been widely used to study the transport of radiation in homogeneous media and in atmosphere:

ATR: R. J. Harris, J. H. Lonergan, and L. Huszar, "Models of Radiation Transport in Air – The ATR Code," SAI-71-557-LJ, La Jolla, CA, Science Applications, Inc., May 1972.

CDR: J. E. Campbell, S. A. Dupree, and M. L. Forsma, "CDR: A Program to Calculate Constant Dose Range from a Point Source of Radiation in the Atmosphere," NWEF 1081, Kirtland AFB, NM: Naval Weapons Effects Facility, 1971.

SMAUG: Harry M. Murthy, "A User Guide to the SMAUG Computer Code," AFWL-TR-72-3, Air Force Weapon Laboratory, Kirtland AFB, NM, May 1972; "A User Guide to the SMAUG-II Computer Code," Air Force Weapon Laboratory, Kirtland AFB, NM, 4 March 1981.

MCNP3B: David L. Monti, "High Altitude Neutral Particle Transport Using the Monte Carlo Simulation Code MCNP with Variable Density Atmosphere," Thesis, Air Force Institute of Technology, AFIT/GNE/ENP/91M-6, Wright-Patterson Air Force Base, Ohio, March 1991; "Modifications to MCNP, Version 3B: Incorporating a Variable Density Atmosphere," Air Force Institute of Technology Technical Report: AFIT/EN-TR-91-2, Wright-Patterson AFB, OH, March 1991.

Mass Integral Scaling Approach

LA-UR-09-00???

If identical isotropic point sources are placed in two infinite homogeneous air media (medium **A** and medium **B**), and if two slant ranges, R_A and R_B , in the two media are related by

$$\rho_A R_A = \rho_B R_B, \quad (1)$$

where ρ is the density in the respective media, then:

$$4\pi R_A^2 F(R_A, E) = 4\pi R_B^2 F(R_B, E). \quad (2)$$

If medium **B** is real air (not homogeneous at all) described by the U. S. Standard Atmosphere, then its “mass range” is defined as $\int \rho_B(Z) dR_B$ and Eq. (1) can be written as:

$$\rho_A R_A = \int \rho_B(Z) dR_B \quad (3)$$

Deterministic Adjoint Weight Window Generator (DAWWG) A Variance Reduction Technique

Motivation

- Research (Wagner, et. al., Larsen, et. al.) have shown that deterministic methods have been effective in generating weight window lower bounds. The methodology is mature.
- These method(s) can be automated, saving user set up time.
- Less user experience with tuning variance reduction techniques is required.

Overview

- The MCNP geometry is “converted” to a mesh representation.
 - Combinatorial geometry is “sampled” based on the weight window mesh.
 - Unstructured mesh can be either “sampled” or used directly.
- A computationally cheap adjoint problem is solved using an S_n deterministic particle transport or diffusion code.
 - Deterministic solution only needs to be “good enough”.
- The weight window lower bounds are calculated from the adjoint flux.
- Source biasing is set so that particles are produced with weights within the weight windows.

Chronology For MCNP

- **Circa 2005: Jeremy Sweezy's initial implementation.**
 - script driven
 - PARTISN deterministic solver used in stand alone mode.
- **Summer 2009: Travis Trahan (summer student) resurrects pervious work and starts move away from scripts.**
- **Summer 2009: funding secured for FY10 to upgrade standalone Capsaicin (S_N , unstructured mesh) for 3-D neutronics & adjoint capability, add collision importance to PARTISN, replace AVATAR capability in PARTISN with one that relies on binary files.**
- **Beyond FY 2010: finish linking to deterministic codes & implementing methodology in MCNP, testing, documentation, publication.**

Concerns

- **Multi-group cross sections are “limited”.**
 - Funding needed to “produce” more cross section sets in NDI format
- **Work may be needed to “fix up” unstructured mesh geometries for deterministic code use or to make deterministic solvers more forgiving with gaps and overlaps.**
- **Others?**